

2 **The Continuing Quest for High T_c –**
3 **Superconductors: A Review**

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8
9 **ABSTRACT**

Superconductors, perhaps a mirage in the eyes of most chemists for many years, was pushed to the forefront of chemical interest with the synthesis and characterization of the first “high - temperature” superconductor (Sr-doped La₂CuO₄ or 2-1-4) in 1986. The first superconductor with a T_c above the boiling point of liquid nitrogen (YBa₂Cu₃O₇ which is often referred to as “1-2-3” compound because of Y: Ba: Cu stoichiometry) was synthesized a year later. This was a time in which superconducting materials created unprecedented excitement in scientific community; it was a year in which the public was captivated by new vistas in technology. The title of the article reflects that there is a crave for superconducting materials. This paper attempts to elucidate the technological advantage of superconductors and reflects the notation that chemistry can have a positive impact upon our lives. Firstly, the historical background of superconductivity is presented then the theory of superconductivity and the description of physical and chemical principles upon which it rests is given. Finally, some prospects for the future applications of these new materials are discussed.

10 *Keywords: High T_c superconductor, 1, 2, 3, compound; liquid nitrogen; Y:Ba:Cu; Phonons*

11
12 **1. INTRODUCTION**

13 Superconductivity is the ability of certain metals, alloys and ceramic materials to let electrical
14 current flow with no electrical resistance and energy dissipation [1]. A superconductor is
15 generally considered as a conventional superconductor if it can be explained by Bardeen,
16 Cooper and Schreiffner (BCS) theory. Conventional superconductors can be either type-I or
17 type-II. Most of the elemental superconductors are conventional except Niobium and
18 Vanadium which are type-II. While other elemental superconductors are type-I but the
19 recently advanced high temperature superconductors are type II or unconventional [1].

20 Since the initial discovery of superconductivity in 1911 by KamerlinghOnnes [2], great
21 advances have been made in both the understanding of the phenomenon and the materials
22 which exhibit it. The materials which exhibit superconductivity have steadily grown in number
23 and variety, but the number used in practical, commercial applications is still rather small.
24 Even though few theories have been developed and widely accepted, extensive researches
25 are underway regarding the deep understanding of the concept of superconductivity, the
26 chemical composition of different superconductors, and to use them as a substitute of the
27 normal conductors in the commercial world. But Understanding the mechanism of high-
28 temperature super-conductivity has remained a subject of much interest since its
29 experimental discovery in the cuprates in 1986 [3].

30 The most important characteristic of any superconductor, from the viewpoint of practical
31 applications, is the maximum electrical transport current density that the superconductor is
32 able to maintain without resistance. This statement is equally true for large-scale
33 applications, such as power transmission lines, electromagnets, transformers, fault-current
34 limiters and rotating machines, as well as for small-scale electronic applications [4].

35 Therefore in this review, the events in the historical developments of superconductivity
36 including the theories, physical and chemical properties which describe the nature of
37 superconductors, the recent advancement and challenges in the application of high
38 temperature super conductors (HTSC) and the future perspectives of high temperature
39 superconductors will be discussed in detail.

40 2. HISTORY OF SUPERCONDUCTIVITY

41 The history of superconductivity as a phenomenon is very rich, consisting of many events
42 and discoveries. The first phenomenon of superconductivity was discovered in 1911 by the
43 Dutch physicist H. KamerlinghOnnes and his assistant Gilles Holst in Leiden. They found
44 that dc resistivity of mercury suddenly drops to zero below 4.2K [5]. In their subsequent
45 experiment they cooled a mercury sample to 3K and again observed that the resistance was
46 “practically zero”. As the temperature of the sample was slowly raised, near 4K, the
47 resistance abruptly increased, by more than three orders of magnitude over a temperature
48 interval much less than 0.1K.

49 Soon after two decades of the discovery of superconductivity, in 1933, W. Meissner and R.
50 Ochsenfeld discovered in Berlin one of the most fundamental properties of superconductors:
51 perfect diamagnetism. The applied magnetic field, below transition temperature, gets to
52 magnetize the substance in an opposite direction, so it shows a negative magnetic
53 susceptibility. They found that the magnetic flux is expelled from the interior of the sample
54 that is cooled below its critical temperature in weak external magnetic fields. Thus, they
55 found that no applied magnetic field is allowed inside a metal when it becomes
56 superconducting. This phenomenon is known today as the Meissner effect [4, 5].

57 Following the discovery of the expulsion of magnetic flux by a superconductor (the Meissner
58 effect) the brothers F. and H. London together proposed in 1935 two equations to govern the
59 microscopic (local) electric and magnetic fields. These two equations provided a description
60 of diamagnetism of superconductors in a weak external field. In the framework of the two-
61 fluid model, the London equations, together with the Maxwell equations describe the
62 behavior of superconducting electrons. The London equations explained not only the
63 Meissner effect, but also provided an expression for the first characteristic length of
64 superconductivity, namely what became known as *the London penetration depth* λ_L [5, 7].

65 Also in 1950, V. Ginzburg and L. Landau proposed an intuitive, phenomenological theory of
66 superconductivity. The equations derived from the theory are highly non-trivial, and their
67 validity was proven later on the basis of the microscopic theory. The Ginzburg-Landau
68 theory played an important role in understanding the physics of the superconducting state.
69 This theory is able to describe the behavior of superconductors (both conventional and
70 unconventional) in strong magnetic fields. The Ginzburg-Landau theory provided the same
71 expression for the penetration depth as the London equations and also an expression for the
72 second characteristic length ξ_{GL} , called the coherence length [5, 7, 8].

73 The first widely-accepted theoretical understanding of superconductivity was advanced in
74 1957 by American physicists John Bardeen, Leon Cooper, and John Schrieffer which is
75 called the “BCS theory”. They described, in great detail, the unique electrical and thermal

76 properties of the superconducting state. The central concept in the BCS theory are the
77 cooper pair of electrons in which such electrons pass through a crystal lattice, the lattice
78 deforms inward towards the electrons generating sound packets known as "phonons". These
79 phonons produce a trough of positive charge in the area of deformation that assists
80 subsequent electrons in passing through the same region in a process known as phonon-
81 mediated coupling. For their pioneering work in developing their theory, they received the
82 Nobel Prize in Physics in 1972 [4, 6, 7].

83 Quantum-mechanical tunneling of Cooper pairs through a thin insulating barrier (of the order
84 of a few nanometers thick) between two superconductors was theoretically predicted by B.
85 D. Josephson in 1962. The DC Josephson effect and AC Josephson effect collectively called
86 the Josephson effects of mathematical predictions were later experimentally approved and
87 played a special role in superconducting applications. [5, 9].

88 During the first 75 years after the discovery of superconductivity in mercury, more than 5,000
89 elements, compounds and alloys were discovered to exhibit superconductivity at a
90 temperatures below about 23 K [6]. But the idea of High-Temperature Superconductivity was
91 not still developed. In 1964 Little hypothesized that high-temperature super- conductivity can
92 be realized in one-dimensional structures on the basis of the exciton mechanism. According
93 to his hypothesis, the system should contain long linear conducting molecules (polyenes or
94 polymers with metallic atoms in an organic matrix), and on each side of these molecules
95 groups of atoms with a high electronic polarizability should be located. These side branches
96 should guarantee the mutual attraction of the conducting electrons as a result of the
97 exchange of molecular excitations of the electronic type [10].

98 The notable events in this regard were the discovery of Nb₃Sn in 1954 and its wire
99 development in 1961. They were followed quickly the discovery of other A15 materials and
100 NbTi, the currently most-used material. These enabled the construction of electromagnets
101 which can produce much higher magnetic fields than conventional copper-wire
102 electromagnets [11].

103 In 1977, Ginzburg and coworkers at the P. N. Lebedev Physical Institute in Moscow
104 published a book entitled "High-Temperature Superconductivity". In this book, they
105 discussed the possibility of high-temperature superconductivity in many types of materials,
106 including quasi-one and- two-dimensional materials. The appearance of this book excited
107 scientists in many countries and the discovery of HTS materials aroused tremendous
108 excitement, because many of them are superconducting and carry significant current above
109 the boiling point of liquid nitrogen at 77.4 K (-196 °C) [10,12].

110 The true history of high T_c superconductors was began in 1986 with the discovery of
111 superconductors on the system Ba-La-Cu-O having the critical temperature 36K by Karl
112 Muller and Johannes Bednorz in IBM research laboratory. They studied the phenomenon of
113 high-temperature superconductivity in the cuprateperovskites materials and they proposed
114 high-T_c superconductivity is achieved when a moderate density of electrons or holes is
115 introduced into the parent antiferromagnetic phases of the cuprate. This opened a new
116 branch of high T_c superconductivity namely "High T_c superconductivity" [13,14]. Soon after
117 this many other oxide based superconductors were discovered having transition
118 temperatures (T_c) up to 138 K (164 K under pressure) [2].

119 In 1987, the groups at the Universities of Alabama and Houston under the direction of M. K.
120 Wu and P. W. Chu, respectively, jointly announced the discovery of the 93 K superconductor
121 Y-Ba-Cu-O. Electron-doped high-T_c cuprates were also first discovered by Tokura et al. in
122 1989 [5]. Pronounced features different from those in the hole-doped cuprates is not only the

123 electron carrier, but also the reduction annealing, which is absolutely essential to obtain the
124 superconductivity in the electron-doped cuprates [15].

125 In 1990, A. S. Davydov presented a theory of high-T_c superconductivity based on the
126 concept of a moderately strong electron-phonon coupling which results in perturbation theory
127 being invalid [5]. The theory utilizes the concept of bisolitons, or electron (or hole) pairs
128 coupled in a singlet state due to local deformation of the -O-Cu-O-Cu- chain in the CuO₂
129 planes.

130 In 2005, a surprising result was obtained in thin films of the electron-doped cuprates with the
131 Nd₂CuO₄-type structure (the so-called T' structure), in which the superconductivity appears
132 without electron doping. Following observations of the undoped (Ce-free) superconductivity
133 in the parent compounds and the suggestion of a new phase diagram opened a new era of
134 research in the high-T_c superconductivity. Yet to date, the mechanism of the undoped
135 superconductivity is unclear [15].

136 137 3. THEORY OF SUPERCONDUCTIVITY

138 Superconductivity is not a universal phenomenon. It shows up in materials in which the
139 electron attraction overcomes the repulsion [5]. Even though the core concept of
140 superconductivity is the above phenomenon, there are different theories regarding the
141 nature, physics of conductivity, type of conductors, physical state of the superconductors etc.

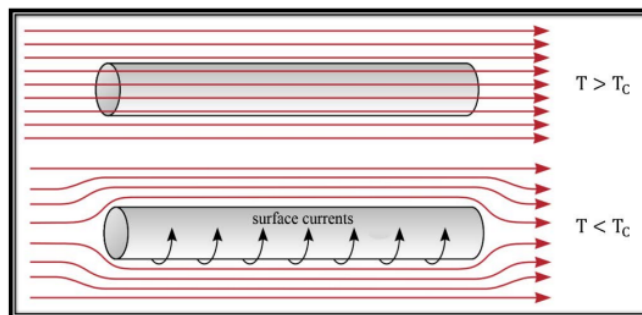
142 3.1. London Theory

143 In 1935 F. London and H. London described the Meissner effect and zero resistivity of the
144 superconducting material by taking $E = 0$ and $B = 0$ in the Maxwell's electromagnetic
145 equation.

$$146 \quad B(x) = B_0 \exp(-x/\lambda_L)$$

147 This first London equation is consequence of the perfect magnetism [16]. According to this,
148 the applied field does not suddenly drop to zero at the surface of the superconductor but
149 decay exponentially according to the equation.

150 According to this theory a magnetic flux penetrating through a superconducting ring or a
151 hollow superconducting cylinder can have values equal to nh/e , where n is an integer
152 (Figure 1). The flux quantization has been confirmed experimentally, but the quantum of flux
153 has been found to be $h/2e$ rather than h/e . This unit of flux is called a fluxoid [14].



154

155 Fig.1. Meissner effect later approved by London theory [36]

156 3.2. Ginzburg - Landau Theory

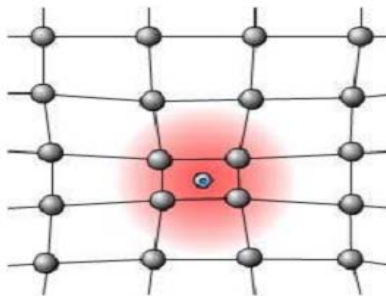
157 In 1950, seven years before the microscopic theory of Bardeen, Cooper and Schrieffer (BCS)
158 was published, Ginzburg and Landau proposed a macroscopic, phenomenological theory of
159 superconductivity to describe the properties of superconductors for temperatures near the
160 critical temperature. There is no concept of the coherence length ξ in the London equation.
161 Ginzburg and Landau proposed the phenomenological theory of the superconductor by
162 using the order parameter $\psi = |\psi_0|\exp(i\theta)$ [4, 16]. According to the Ginzburg–Landau theory,
163 when ψ is small, the Helmholtz free energy per unit volume of the superconductor is
164 expanded as a Taylor series of $|\psi|^2$. Thus we have

$$165 F_s = F_n + \alpha|\psi_0|^2 + (\beta/2)|\psi_0|^4$$

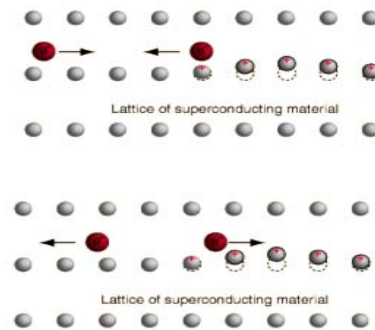
166 Where F_s and F_n are the Helmholtz free energies for the superconducting state and the
167 normal conducting state, respectively, and α and β are the expansion coefficients.

168 3.3. BCS Theory

169 The first widely-accepted theory to explain superconductivity put forth in 1957 by John
170 Bardeen, Leon Cooper, and John Schrieffer. The theory asserts that, as electrons pass
171 through a crystal lattice, the lattice deforms inward towards the electrons generating sound
172 packets known as "phonons". These phonons produce a trough of positive charge in the
173 area of deformation that assists subsequent electrons in passing through the same region in
174 a process known as phonon-mediated coupling (Figure 2). The consequence is an attractive
175 interaction between the two electrons. Now, the correlated pair of electrons is referred to as
176 a Cooper pair. Cooper pairs occupy a collective state and move through the crystal lattice
177 unimpeded [4, 7, 14].



Lattice deformation



178

179 Fig.2. Schematic representation of phonon-electron attraction in a lattice [7]

180 When sufficient energy is supplied to Cooper pairs, they will decouple and the
181 superconducting state ceases. Since the momentum is proportional to current density, the
182 energy required to break Cooper pairs implies the existence of a critical current density J_c
183 [18]. When J_c is exceeded, the Cooper pairs are destroyed. The energy required to break
184 Cooper pairs also comes in the form of thermal energy from the vibrations of the crystal
185 lattice and magnetic energy from magnetic fields. The former implies the existence of a
186 critical temperature T_c , mentioned above, while the latter suggests the existence of a critical
187 magnetic field H_c . The critical field depends on material and is temperature dependent
188 according to the following equation

189
$$H_c = H_0[1 - (T/T_c)]^2$$

190 Where H_0 is the value of the critical field strength as $T \rightarrow 0$.

191 The BCS theory in its original form cannot be applied to the high temperature
192 superconductors because the BCS predicted temperature are too low and because of non
193 occurrence of the BCS predicted isotope effect [7, 16].

194 **3.4. Possible mechanism of conductivity in HTSC**

195 Until Fe-based superconductors were discovered in 2008, the term high-temperature
196 superconductor was used interchangeably with cuprate superconductor for compounds such
197 as bismuth strontium calcium copper oxide (BSCCO) and yttrium barium copper oxide
198 (YBCO) [17].

199 The question of how superconductivity arises in high-temperature superconductors is one of
200 the major unsolved problems of theoretical condensed matter physics as of 2010 [7, 12]. The
201 mechanism that causes the electrons in these crystals to form pairs is not known [5]. Despite
202 intensive research and many promising leads, an explanation has so far eluded scientists.
203 One reason for this is that the materials in question are generally very complex, multi-
204 layered crystals (for example, BSCCO), making theoretical modeling difficult.

205 There have been two representative theories for HTS. Firstly, it has been suggested that the
206 HTS emerges from antiferromagnetic spin fluctuations in a doped system [19]. According to
207 this theory, the pairing wave function of the cuprate HTS should have a $d_x^2 - y^2$ symmetry.
208 Thus, determining whether the pairing wave function has d-wave symmetry is essential to
209 test the spin fluctuation mechanism. That is, if the HTS order parameter (pairing wave
210 function) does not have d-wave symmetry, and then a pairing mechanism related to spin
211 fluctuations can be ruled out. Secondly, there was the interlayer coupling model, according
212 to which a layered structure consisting of BCS-type (s-wave symmetry) superconductors can
213 enhance the superconductivity by itself. By introducing an additional tunneling interaction
214 between each layer, this model successfully explained the anisotropic symmetry of the order
215 parameter as well as the emergence of the HTS [5, 7, 20, 21].
216

217 **4. PROPERTIES OF SUPERCONDUCTING STATES**

218 Among the many superconducting properties some basic and significant physical and
219 chemical properties are discussed. Along with those properties, the superconducting property
220 of cuprates particularly YBCO is presented.

221 **4.1. Critical temperature**

222 The most commonly known property of the superconducting state is that a superconductor is
223 characterized by a critical temperature, below which the material exhibits "zero" resistance
224 [6]. The critical temperature is mainly determined by chemical composition and structure,
225 [22]. Superconducting cuprate materials have higher transition temperatures. $YBa_2Cu_3O_{7-x}$
226 (YBCO) has a T_c of 93 K which is significant because it is greater than the boiling point of
227 liquid nitrogen at atmospheric pressure as shown in Table 1. Bi-Sr-Ca-Cu-O, Tl-Ba-Ca-Cu-
228 O, Hg-Ba-Ca-Cu-O and Hg-Tl-Ba-Ca-Cu-O compounds have higher critical temperatures as
229 shown in the table [23]. All these high temperature superconductors have highly anisotropic
230 crystal structures, containing layered CuO_2 planes in which the superconducting charge
231 carriers are thought to be localized.

232

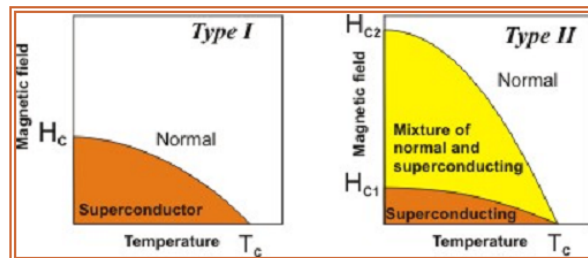
Table.1. Some high Tc cuprate based superconductors [23]

Compounds	Tc K	Compounds	Tc K
YBA ₂ CU ₃ O ₇	93	TL ₂ CABA ₂ CU ₂ O ₈	119
YBA ₂ CU ₄ O ₈	80	TL ₂ CABA ₂ CU ₂ O ₁₀	128
Y ₂ BA ₄ CU ₇ O ₁₅	93	TLCABA ₂ CU ₂ O ₇	103
BI ₂ CASR ₂ CU ₂ O ₈	92	TLCA ₂ BA ₂ CU ₃ O ₈	110
BI ₂ CA ₂ SR ₂ CU ₃ O ₁₀	110	TL _{0.5} PB _{0.5} CA ₂ SR ₂ CU ₃ O ₉	120
HGBA ₂ CA ₂ CU ₃ O ₈	135	HG _{0.8} TL _{0.2} BA ₂ CA ₂ CU ₃ O _{8.33}	138

233

234 **4.2. Critical magnetic fields**

235 Superconducting state can be destroyed by a sufficiently strong magnetic fields. But for type-
 236 II superconductor, there are two critical fields, the lower critical field H_{c1} and the upper critical
 237 field H_{c2}. In applied fields less than H_{c1}, the superconductor completely expels the field, just
 238 as a type-I superconductor does below H_c. At fields just above H_{c1}, flux, however, begins to
 239 penetrate the superconductor. Type II superconductors are the most technologically useful
 240 because the second critical field can be quite high, enabling high field electromagnets to be
 241 made out of superconducting wire as shown in Figure 3. Wires made from say niobium-tin
 242 (Nb₃Sn) have a H_{c2} as high as 24.5 Tesla – in practice it is lower. This makes them useful for
 243 applications requiring high magnetic fields, such as Magnetic Resonance Imaging (MRI)
 244 machines, [5].



245

246

247

Fig. 3. Schematic plot of magnetic field vs temperature for type-I and type-II superconductors.[36]

248

4.3. Critical current

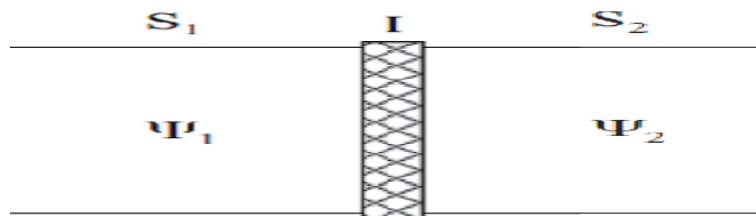
249 Another important characteristic of a superconductor is the maximum transport current which
250 can flow without dissipation. This is called the critical current, I_c . Its value is very sample
251 dependent and can be affected by the sample shape and material quality. There is a
252 criterion that says a superconductor loses its zero resistance when, at any point on the
253 surface, the total magnetic field strength, due to the transport current and applied magnetic
254 field, exceeds the critical field strength H_c . This quantity, I_c , is called the critical current.
255 I_c depends on the external magnetic field experienced by the superconductor and has typical
256 values of the order of 10^6 - 10^8 A cm^{-2} depending on the sample temperature [24].

257 4.4. Zero resistance

258 Every superconductor has zero resistivity, i.e. infinite conductivity, for a small-amplitude dc
259 current at any temperature below T_c . This property of the superconducting state was
260 demonstrated by inducing a small-amplitude dc current around a closed ring of a
261 conventional superconductor. The experiment continued over two and a half years—there
262 was no measurable decay of the current. This means that the resistivity of a superconductor
263 is smaller than 10^{-24} Ωm . This value is 18 orders of magnitude smaller than the resistivity of
264 copper (1.7×10^{-6}) at room temperature. Such a value of resistivity in a superconductor
265 implies that the current lifetime in a super-conducting ring in zero magnetic field is not less
266 than 10^5 years [5, 24].

267 4.5. Josephson Effects

268 As with semiconductor and thermocouple devices, there is the concept of a junction in a
269 superconductor when it joins with an insulator. Two superconductors are separated by a thin
270 insulating layer whose thickness is so small that Cooper pairs can pass through by the
271 tunneling effect as given in Figure 4. This geometry is called the Josephson junction.



272
273 Fig. 4. Structure of Josephson junction (where s_1 and s_2 are two kinds of
274 Superconductors, I is thin insulating layer) [36].

275 The Josephson Effect is the phenomenon of super current (A current that flows indefinitely
276 long without any voltage applied) across a device which is known as Josephson junction.
277 The Josephson Effect is an important basis of superconducting electronics applications and
278 has widespread applications in many instruments such as voltage reference,
279 superconducting cavities, superconducting filters and Superconducting Quantum
280 Interference Devices (SQUIDs) [1].

281 4.6. Thermal property

282 4.6.1 Entropy

283 A marked decrease in entropy is observed during normal to superconductivity transition
284 near the critical temperature; which indicates that the superconducting state is more ordered
285 than normal state.

286 **4.6.2 Specific heat**

287 From the specific heat study of superconductor we can get the information about the
288 existence of the band gap in superconductors. As we know specific heat of normal metal,

289
$$C_n = \gamma T + \beta T^3$$

290 Where, γT =specific heat term

291 βT^3 =contribution of lattice vibration at low temp

292 Specific heat of superconductor shows a jump at T_c since the superconductivity affects
293 electron mainly. So, the lattice vibration part remains unaffected. By this substitution the
294 electronic specific heat C_{es} shows the exponential curve.

295
$$C_{es}(T) = A \exp(-\Delta_{TK} / k_B T)$$

296 And this indicates the existence of finite gap in the superconductor [14].

297 **4.7. Superconducting Energy Gap**

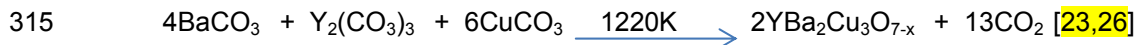
298 According to the BCS theory, the Cooper pairs are bound together with an energy E given by
299 the relationship

300
$$E = h\nu = 3.52 k_B T_c$$

301 Where h is Planck's constant ($h = 6.6262 \times 10^{-34}$ Joule-sec), ν is the frequency, k_B is Boltzmann
302 constant ($k_B = 1.3806 \times 10^{-23}$ Joules/Kelvin) and T_c is the superconducting transition
303 temperature. The photon frequency corresponding to the energy gap of a superconductor
304 with a $T_c = 1$ K is about 73 GHz. Thus the energy gap of most superconductors corresponds
305 to photons in the terahertz or far infrared frequency regions of the spectrum [6]. Absorption
306 of energy in that region may cause promotion of electron from super conducting state to
307 metallic state. Additional evidence for reduced free energy in super conducting state comes
308 from heat capacity measurements. These show the superconducting state to have lower
309 entropy and hence to be more ordered than the metallic state [25].

310 **5. SUPERCONDUCTIVITY IN CUPRATE OXIDES**

311 Since 1987, cuprate superconductors with $T_c > 77K$ have been the centre of intense
312 interest. One of the first to be discovered was $YBa_2Cu_3O_7$ made by reaction given below.
313 The oxygen content of the final material depends on reaction conditions (e.g. temperature
314 and pressure).



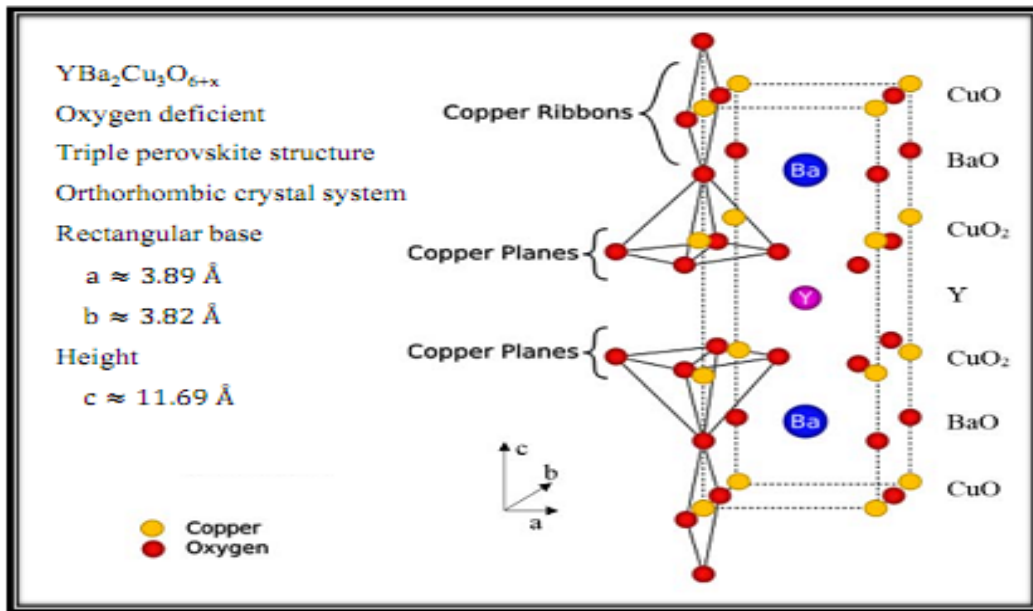
316 It is related to the perovskite structure as follows: by tripling the perovskite (ABO_3) unit
317 cell and substituting one yttrium atom for every third barium atom, the formula $Y_1Ba_2Cu_3O_9$
318 results. However, a little more than two oxygen vacancies are required for
319 superconductivity there are position of nine atoms of O in the unit cell but only 7 are occupied
320 [27,28].

321 There are systematic oxygen atom vacancies in the unit cell compared to stack of simple
322 perovskite unit cells. These occur between adjacent copper atoms along the c axis. The
323 vacancies are in the yttrium atom plan. There are also vacancies between copper atoms
324 along the a axis in the copper-and-oxygen planes that lie between the planes of barium

325 atoms [26]. This non-stoichiometry is denoted by the x in the chemical formula $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$.
 326 When $x = 1$, the O(1) sites in the Cu(1) layer are vacant and the structure is tetragonal. The
 327 tetragonal form of YBCO is insulating and does not superconduct. Increasing the oxygen
 328 content slightly causes more of the O(1) sites to become occupied. For $x < 0.65$, Cu-O chains
 329 along the b axis of the crystal are formed. Elongation of the b axis changes the structure to
 330 orthorhombic, with lattice parameters of $a = 3.82$, $b = 3.89$, and $c = 11.68$ Å. Optimum
 331 superconducting properties occur when $x \approx 0.07$, i.e., almost all of the O(1) sites are
 332 occupied, with few vacancies.

333 In experiments where other elements are substituted on the Cu and Ba sites, evidence has
 334 shown that conduction occurs in the Cu(2)O planes while the Cu(1)O(1) chains act as
 335 charge reservoirs, which provide carriers to the CuO planes. However, this model fails to
 336 address superconductivity in the homologue Pr123 (praseodymium instead of yttrium). This
 337 (conduction in the copper planes) confines conductivity to the a - b planes and a large
 338 anisotropy in transport properties is observed (Figure 5). Along the c axis, normal
 339 conductivity is 10 times smaller than in the a - b plane. For other cuprates in the same general
 340 class, the anisotropy is even greater and inter-plane transport is highly restricted [29].

341 The oxidation state of copper in $\text{Y}_1\text{Ba}_2\text{Cu}_3\text{O}_7$ is unusual. If we assume that Y, Ba and O have
 342 their usual oxidation state +3, +2 and -2 respectively, for charge balance, the copper must
 343 have an average of +2.33. This may be rationalized in terms of one third of the copper is Cu^{3+}
 344 and the remainder is Cu^{2+} [1,23,25,26,28].



345

346 Fig. 5. The crystal structure of $\text{YBa}_2\text{Cu}_3\text{O}_7$ oxygen deficient compound [36]

347 6. APPLICATION OF SUPERCONDUCTORS

348 Due to the unique properties of high temperature superconducting materials, their
 349 applications are becoming attractive continuously with the improvement in properties of
 350 superconductors. The search for applications has always been a driving force for
 351 superconductor material science. Right from the discovery, it had been envisioned that

352 superconducting coils with high persistent current might be used to produce strong magnetic
353 field.

354 In the last 20 years, many applications of high temperature superconductors have been
355 developed in very diverse fields. Applications of superconductivity can be found in
356 transportation (maglev trains), marine and military (propulsion motors, degaussing systems
357 and EMP weapons), particle research (large hadron collider, proton-antiproton collider and
358 electron proton collider etc.), power generation and distribution (fault current limiters,
359 superconducting wires, superconducting magnetic energy storage systems and
360 superconducting transformers etc.), information technology & computing (quantum
361 computers, quantum cryptography and high performance computers etc), electronics &
362 telecommunications (Superconducting Quantum Interference Device (SQUID), single-
363 quantum flux devices (SFQ devices), and cellular filters etc.) and medical diagnostic systems
364 (magnetic resonance imaging (MRI))[30]. In this section some typical applications of high
365 temperature superconductors are discussed.

366 The fundamental technologies for applications of superconductivity are discussed here with
367 the classifications of Superconducting bulks, Superconducting tapes and Superconducting
368 devices.

369 **6.1. Superconducting bulk and its applications**

370 Materials for superconducting bulk are $REBa_2Cu_3O_7$ where RE is Sm, Nd, Gd or Y. The bulk
371 is made using the oxygen controlled melt growth (OCMG) method under a low partial
372 pressure of oxygen except for YBCO where the quenched melt growth (QMG) method is
373 used, here the superconducting bulk YBCO is grown using half-melt materials at a high
374 temperature of nearly 1000 °C, followed by a very slow cooling. The bulk almost has a
375 single-crystal like structure, and a large critical current of more than $10^5 A/cm^2$ is obtained at
376 liquid-nitrogen temperature [12,31].

377 The special feature of the bulk is that the introduction of pinning centres is easily carried out
378 by controlling the density of fine Y_2BaCuO_5 particles as pinning centres in the bulk. The most
379 distinguishing characteristic is that it is possible to trap a strong magnetic field of 2 to 3 T,
380 even at liquid-nitrogen temperature. This value is much higher than the magnetic field of an
381 ordinary permanent magnet. Recently, the mechanical strength of bulk has been increased
382 by polymer impregnation, and this bulk has trapped a very high magnetic field of 17 T at
383 30K. The applications of the bulk are of two types:

384 (i). As the pinning force of the magnetic flux is so strong, the outside magnetic field cannot
385 penetrate the bulk in the superconducting state. This results in strong levitation force of
386 about 15 Kg/cm² when the bulk is close to an ordinary permanent magnet at 77 K. This
387 phenomenon is exploited to make Flywheel electric storage systems that can store about 10
388 KW/h, based on frictionless superconductor bearings. Using bulk HTS self-centering
389 bearings allows levitation and rotation in a vacuum, thereby reducing friction losses.
390 Conventional flywheels suffer energy losses of 3-5% per hour, whereas HTS based
391 flywheels operate at <0.1% loss per hour. Large and small demonstration units are in
392 operation and development [1,12,31].

393 (ii). When the bulk is in normal state, a magnetic field applied from outside is uniformly
394 distributed throughout the bulk. But after the bulk is cooled to below the critical temperature
395 the magnetic field is quantized and quantized flux is pinned by strong pinning centres. Then
396 when the external field is removed the quantized flux is left inside and behaves like a
397 permanent magnet. One fruit full application of this bulk is in water cleaning using the
398 magnetic separation effect. The impurity particles in water join magnetic particles, and they

399 are removed from the filter by the strong magnetic field of the bulk. The operation can be
400 continuous and results have been impressive which is 100 times efficient than the present
401 cleaning system available today. Other opportunities are in materials manufacturing such as
402 semiconductor production in high magnetic field and induction heating [1, 12, 31, 32].

403 **6.2. Superconducting tapes**

404 **6.2.1. First-generation superconducting tapes**

405 The first generation of superconducting tape using high temperature superconductors is the
406 so-called Silver-sheathed Bi-compound Tape. The Bi-compound usually used is
407 $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_{10}$ (Bi-2223). The tape is 4mm in width and 0.25mm in thickness, and inside
408 the sheath there are 55 filaments of Bi-2223 superconductor. Recently, Sumitomo has
409 succeeded in producing higher-quality tapes by introducing a heat-treatment process under
410 a high pressure of 300 atmospheric pressures, and has obtained a critical current of 200A at
411 77 K [12,31,33].

412 **6.2.2. Second-generation superconducting tapes**

413 Although mechanical properties of the first generation wires are reasonably robust and the
414 critical current values have reached more than 10^4 A/cm^2 at 77K and at commercial lengths,
415 there are several weak points in this system in that high critical current density decreases
416 rapidly with increasing magnetic field at liquid nitrogen temperature. Therefore, recent
417 developments in the second generation coated conductors to improve the magnetic field
418 dependence, where a thin film of $(\text{RE})\text{Ba}_2\text{Cu}_3\text{O}_y$ superconductors is deposited on an
419 appropriate substrate with specifications required by the specific application, gain a lot of
420 interest [33]. YBCO tape consists of a thin metal plate, a buffer layer and a superconducting
421 layer has developed with critical current density of $3 \times 10^6 \text{ A/cm}^2$ in a magnetic field of 20 T
422 at approximately 20 K that have law dependence in magnetic field [1, 12, 31].

423 **6.2.3. Electric Power**

424 ***6.2.3.1 Superconducting electric power cable***

425 A possibility of electric energy transportation without loss is the most attractive property of
426 superconductors. A high transition temperature of the cuprate HTS has enabled their
427 practical applications by development of HTS power cables which are able to carry up to ten
428 times greater current and power capacity per cross section than conventional copper cables.
429 HTS coated conductor tapes are used in fabricating power cables. American
430 Superconductor Company (AMSC) produces a single phase power cable using GdBCO
431 tapes that are fabricated by using pulsed laser deposit technique and having greater current
432 density than YBCO tapes. To reduce HTS cable cost three phase power cables were
433 developed by AMSC, Where three 2G wires are concentrically placed around the common
434 central core surrounded by a copper shield. The Triax HTS cable was developed jointly by
435 Southwire and nkt cables companies. This compact construction has several advantages
436 over the single-phase cable: in the three-phase cable, only about one-half of the quantity of
437 HTS wire is needed and the cables cold surface area is reduced, which lower the cost
438 associated with cryogenic cooling equipment. A single HTS Triax cable operating at 13 kV
439 carries 3,000 A, which is equivalent of 18 conventional underground cables [1,32, 34].

440 ***6.2.3.2 Fault Current Limiters***

441 As new generators are added to the network, many local grids face a rising risk of
442 unacceptably high power surges that result from “faults” or short circuits. A fault current
443 occurs in the event of a short circuit caused by lightning, accidental contact between the
444 lines or the ground, etc. In this case, the power current flowing through a local network can
445 increase enormously damaging electrical equipment. To protect the transmission or
446 distribution of electric systems from outages caused by fault currents a special device – a
447 fault current limiter (FCL) is installed in the transmission grid. Conventional line reactors
448 widely used as FCLs have high AC losses and can produce voltage drop in the grid in the
449 case of a fault current. HTS technology enables a new solution: compact, “smart” fault
450 current limiters (FCLs) that operate, passively and automatically, as power “safety valves” to
451 ensure system reliability when individual circuits are disrupted. Taking advantage of the
452 inherent properties of superconductors, they sense such dangerous over currents and
453 reduce them to safe levels by changing state instantaneously, from “super” conductors to
454 resistors when the electric current exceeds the materials critical current.

455 The HTS FCL represents a coated conductor consisted of layers of HTS material within
456 layers of resistive materials. Under normal operating conditions, the current in the cable
457 flows through the HTS layers in the FCL. In the case of a fault, the current exceeds the HTS
458 material’s critical current and the HTS layers become normal. In that case, the current is
459 automatically shunted within a millisecond to flow through the higher resistance layers,
460 effectively quenching the fault current amplitude. The very rapidly operated HTS FCLs
461 greatly reduce damage to electrical equipment caused by system faults. They are fail-safe
462 since they require no external sensing of the current to initiate the transition [1, 32].

463 **6.2.4. Transportation**

464 ***6.2.4.1 Superconducting Magnet for Maglev Trains***

465 Expulsion of magnetic field by superconductors in superconducting state is useful in
466 magnetically levitated trains, called as Maglev trains, which are operational in a few
467 countries. The magnet for maglev train has been made by using Bi-2223 superconducting
468 tape. This magnet has a racetrack shape and consists of twelve pancake coils, of 1m length
469 and 50 cm height. This magnet is operated at 20k and generates a magnetic field of 2.5 T
470 at the center of the magnet. The huge magnetic field that can be sustained by the
471 superconductors is used to levitate and propel the trains. Furthermore, the rate of decay of
472 the persistent current is only 0.5% a day. This magnet is very successful and recently a train
473 using this magnet reached a speed of 500 km/h [1, 35].

474 **6.2.5. Marine**

475 ***6.2.5.1 Superconducting Motor for Ships***

476 The ship propulsion system is undergoing a revolution. In the new system, propellers and
477 the electric motor are directly connected and they are outside the body of the ship and this
478 propulsion system is called a “pod motor”. By employing such a system, freedom of boat
479 design is very much improved, and as a result, energy saving becomes possible.

480 The large size and heavy weight of conventional copper-based electrical propulsion motors
481 and generators has been a barrier to broad adoption of electric propulsion. A HTS motor is
482 most suitable for this pod motor system, because the superconducting motor generates a
483 large torque even at a slow rotating propeller speed of about 100 rpm; furthermore, it is
484 much smaller and much lighter than ordinary motors using Medical Imaging and Diagnostics
485 copper wire. In the United States, various types of superconducting motors are being

486 developed; all of them use Bi-2223 superconducting tape. In Japan, a superconducting
487 motor using YBCO tape has been recently developed, and this is the first such motor in the
488 world [1,32].

489 **6.2.6. Medical application**

490 ***6.2.6.1 Magnetic Resonance Imaging (MRI)***

491 One of the largest commercial power applications of superconductors at present is magnetic
492 resonance imaging (MRI), with thousands of units in hospitals and global sales of several
493 billion US dollars per year. The superconducting portion consists of a “basic” solenoid, which
494 creates the background strong magnetic field that forces hydrogen atoms that exist in the
495 body’s water and fat molecules to accept energy from the magnetic field. These species then
496 release this energy at a certain frequency which can be detected and displayed in the form
497 of an image by a computer [1,2,35].

498 **6.3. Superconducting electronic devices**

499 One of the most prominent phenomena of superconductors is the Josephson tunnelling
500 effect and most of the applications of superconductors in electronics and sensors are based
501 on this phenomenon. Applications of superconductors have arisen in areas other than
502 electromagnets, more specifically in electronics and sensors, wherein the materials required
503 are often small thin films which are much easier to produce in highly perfect forms,
504 especially from HTS. [2].

505 **6.3.1 Superconducting Quantum Interference Devices (SQUID)**

506 SQUIDs are superconducting loops with integrated Josephson junctions which can be used
507 as the most sensitive measurement for magnetic fields, voltage, and related electromagnetic
508 quantities based on the Josephson tunnelling phenomena. The magnetic field resolution of
509 HTS SQUIDs operating at 77 K temperature is about 10 times lower than the commercial
510 LTS SQUIDs operating at 4K [1]. However, a large commercial impact is only expected for
511 HTS SQUID systems that are able to observe magnetic signals even in the presence of
512 disturbing background fields without the burden of magnetic shielding [36]. SQUID
513 magnetometers may be the most sensitive measurement device known. The threshold for
514 SQUID is of the order of 1 fT,(100 billion times smaller than the Earth’s field), making it
515 capable of measuring extremely feeble magnetic fields. Because of their extreme sensitivity,
516 SQUIDs have established themselves as very accurate devices for both
517 Magnetocardiography and Magnetoencephalography. The application of high temperature
518 SQUIDs to a magnetocardiography has been made recently by Hitachi, Ltd. It consists of 51
519 SQUIDs on one plate and is used for the diagnosis of human heart diseases [1].

520 **6.3.2 Microwave filters**

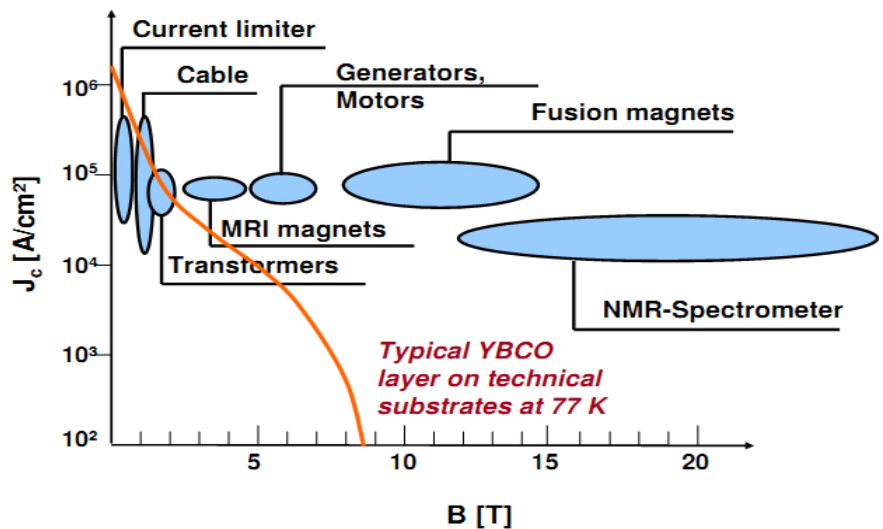
521 The simplest applications of high-quality high-temperature superconducting films with
522 commercial availability are HTS passive RF and microwave filters for wide-band
523 communications and radar. These are based on conventional microstrip and cavity designs
524 with superconductors used for the microstrips or to line the metal cavity. They have the
525 advantages of very low noise and much higher selectivity and efficiency than conventional
526 filters [12]. In early stages of development it was hoped that many filtering systems would be
527 used in the base stations of portable telephone systems. However, at present, only several
528 thousand filters are employed in the world, because of the cost of the cooling system [10].

529 **6.3.3 Single-Quantum Flux (SFQ) Devices**

530 Another important application of **High Temperature Super Conductors**, HTSC is the single
531 quantum flux device (SFQ). The principle operation of SFQ device is that in a
532 superconductor ring of a SQUID the magnetic field is quantized, and by applying a current
533 pulse to the ring, the Josephson junction reaches a normal state for a short duration, and the
534 quantized flux appears or disappears in the ring depending on the original state. The state of
535 the flux in the ring is 0 or 1, and responds to the 0 or 1 of an information signal and so logical
536 circuits can be made by combining SFQ devices [10,31]. The electricity consumption of the
537 SFQ circuits is very small, 0.1 mW per one logic gate, and this is about one hundredth that
538 of semiconductor circuits. The operation speed is about 100 GHz, and this is one hundred
539 times faster than that of semiconductor circuits [1].

540 In the past fifteen years, the circuit technology of SFQ circuits using Nb-based SFQs has
541 made a good progress, the integration of more than 10^4 junctions has become possible,
542 high-speed shift registers and high-speed switching systems have been developed. Special
543 feature of SFQ circuits based on high-temperature superconductors are a high-speed
544 operation of over 100 GHz and a high-temperature operation at approximately 40 K. On the
545 other hand, the integration is more difficult than that of Nb-based SFQ circuits owing to the
546 complexity of process technology (Figure 6). A high-speed sampler and a toggle flip-flop
547 circuit are examples of SFQ circuits with about twenty Josephson junctions and operating at
548 about 40 K. The high-speed sampler is very useful for observing waveforms of a very short
549 pulse greater than 40 GHz, which will be popular in future, communication systems. And the
550 toggle flip-flop circuit generates two output pulses for one input, and can be used as
551 demultiplexer and this circuit is operates faster than 370 GHz [10]. At present, the integration
552 of a high-temperature SFQ circuit is limited to about 100 junctions, but it is hoped that this
553 will reach to more than 500 junctions.

554



555

556 **Fig.6. Required properties for different power applications of HTSC[36].**

556

557 **6.4 Future perspective of HTS**

557

558 The last ~30 years have seen the development of a number of prototypes of new high
559 temperature superconducting power and electronic equipment's that are mentioned above
560 and many others. But, so far only few of them are successfully commercialized. The problem
561 has usually been the cost of the HTS version versus other LTS and non-superconducting
562 solutions. Superconducting motors and generators are one strong example; a number of
563 prototypes have been successfully built and tested to full function, but the economics have
564 not allowed commercialization [2]. High temperature superconductivity does not necessarily
565 require correlated electron systems with complex competing or coexisting orders. Instead, it
566 may be achieved in a phonon-mediated classical superconductor having a high Debye
567 temperature and large electronic density of states at the Fermi level in a material with light
568 atoms and strong covalent bonds [37-39].

569 The future of HTS applications and commercialization is a persistent question and can only
570 be answered once the scientific and technological challenges hindering the large-scale
571 application of HTS are addressed. First, the geometrical shapes of HTS wire are limited. It is
572 mostly available in tape format. Fabrication of HTS wire (tape) requires an expensive and
573 special technology and process machinery.

574 The second challenge has to deal with the material properties that result in high values of TC
575 and HC. These properties include high normal state resistance, proximity of
576 superconductivity to competing Anti-Ferromagnetic (AF) state, large anisotropic ratios and
577 sensitivity of superconducting properties to local nonstoichiometry. These lead to the lack of
578 widespread applications of HTS. In addition, absence of any validated microscopic theory of
579 superconductivity in HTS materials means we cannot predict new superconducting materials
580 with higher TC values.

581 If these challenges of HTS are addressed, then the HTS will bring the features of save
582 energy, compact, lightweight, high performance and save resources, and could play a trump
583 on forthcoming low-carbon society in our life.

584 **7. CONCLUSION**

585 One of the most brilliant events in the 20th century was the discovery of ceramic-type HTS
586 materials that led to extensive research for new discoveries and theories. The physics of
587 HTS is complicated, that the theoretical basis is uncertain yet, and it may be in the far front
588 of solid state physics. Till this moment high temperature superconducting properties has
589 been discovered with in many types of materials and research will continue to synthesize
590 new materials that may expose these new phases in an experimental sense or improve on
591 those already known. We have studied historic view of HTS, theories of superconductivity,
592 and characteristics and different properties of HTS materials. We also studied the application
593 of HTS in different sectors. As the properties of HTS continue to improve, more applications
594 are expected to become a commercial reality. Steady improvement of the HTS materials
595 basis will surely widen this spectrum of applications within near future.

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